Unorthodox (heretic?) approaches for (electro)magnetic fields for compact objects

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Driving goal

- Understand both gravitational and electromagnetic wave emissions from key systems
 - Binary black holes interacting with surrounding media
 - Binary neutron stars
 - Black hole neutron star binaries
- Connect with efforts to understand emissions from single black hole systems
 - E.g. single black hole + disk
 - AGNs, GRBs, etc.

Peculiarities

- Deal with spacetime curvature
 Einstein equations
- Deal with fields describing fluids and electromagnetic fields
 - perfect fluid for matter
 - Electromagnetic fields? Not a unique prescription
 - Far away, EM fields in vacuum
 - Regions with $B^2 >> P \rightarrow$ 'force-free' regime
 - Regions with P >> $B^2 \rightarrow$ inertia of matter dominates, ideal MHD regime

Work towards all 3 regimes in one ongoing (Palenzuela's talk), in the mean time, want to extract valuable information for physics and longer term goals.

Peculiarities contd

• GR related

- Broadly speaking nothing is conserved (beyond ``constraints'')
 - Efforts to truly conserve quantities key in Newtonian (or fixed background) settings while might help are neithre required nor necessarily well defined...
 - Speed of light time step constraint.
 - For possible wide applications (physics, regions, new ideas) would like flexible methods not constraining algorithmic options [e.g. convergence and consistency only requirements]
 - Ex: DivB=0 and E.B=0 constraints → treat both at equal footings.
 For instance, through ``Lagrange multipliers''

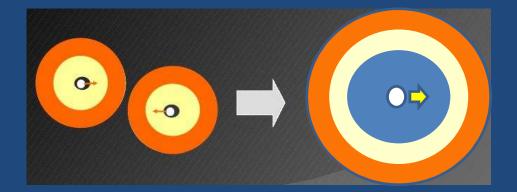
Outline. (Piece-meal approach to EM effects)

- What can curvature do?
 Binary black holes as field stirrers
- What can curvature, and inertia of matter do?
 Binary neutron stars
- Curvature, inertia of matter and horizon
 Black hole neutron stars

In all cases: Where the wild things are (i.e. fall apart)

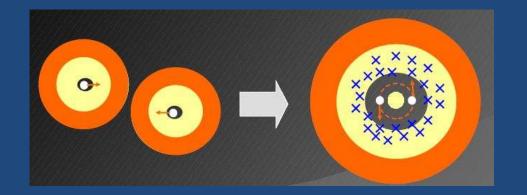
Binary black holes and emissions

- Different possible options.
 - Postmerger events from circumbinary disks around BHs



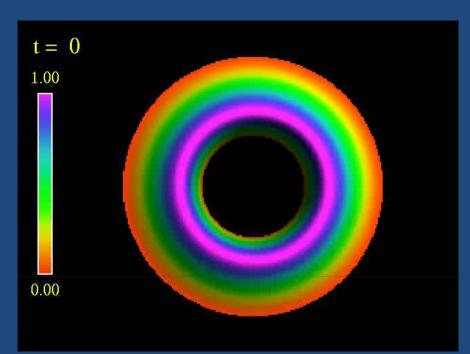
[Milosavljevic-Phinney; Lipai-Loeb; Lipai et.al, Bonning et.al; Bode et.al; Megevand et.al]

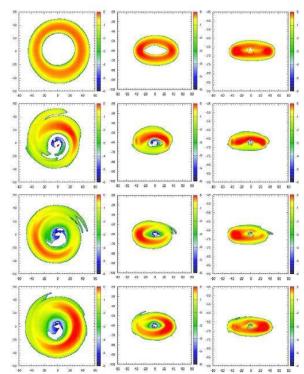
Pre/merger events from gas in between BHs / torques on dis



[Armitage et.al; MacFadyen et.al.; Dotti et.al; Chang. et.al.; Palenzuela et.al.]

Postmerger emission





[Megevand,Anderson,LL,Neilsen;2010] [Megevand, Anderson,Frank,LL,Liebling,Neilsen,Motl; 2009]

Key qn...

• How are possible jets affected?

[eg. Hawley, McKinney, Gammie, Stone, Fragile, Komissarov, etc]. Work in progress [Mackinney, LL]

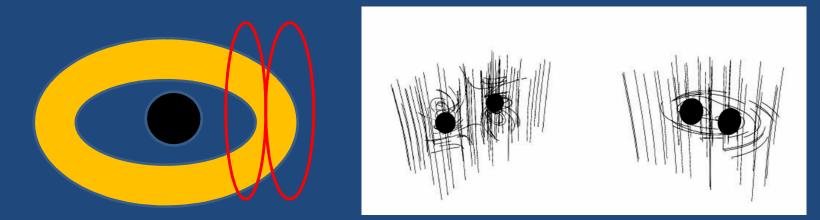
Binary black holes as blenders: 'B-Z' for binaries

How does the curvature influence EM fields?

Ingredients: GR + Maxwell Eqns (ElectroVac)







[Palenzuela, LL, Anderson, Liebling, Neilsen, PRL 2009 ; Palenzuela, LL, Yoshida 2009; Moesta, Palenzuela, Yoshida, Rezzolla, Pollney, LL, submitted]

Approach: Electrovac (GR+EM)

- Einstein equations
- Maxwell equations: now with currents (Force free approx)

Initial setup

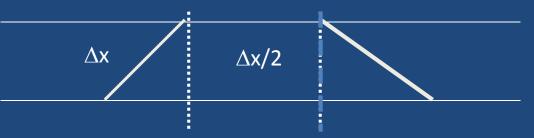
- Quasicircular, equal mass
- Magnetic field as given by a circular loop at far distances ~ constant within computational domain
- Field strength (quite conservative) = 10^4 G
 - For this value, if $M_{\tau}=10^8 M_o$, EM Energy dens ~ 10^{-16} [1/M²]

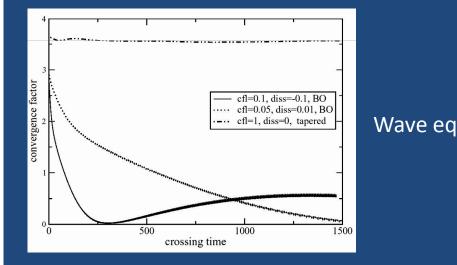
 \rightarrow EM fields won't affect binary dynamics, but the other way around

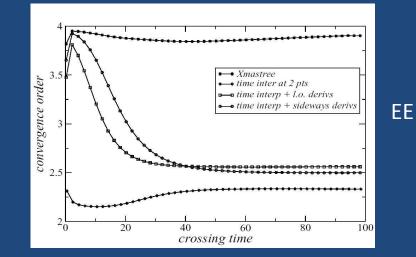
→ results hold for fields ~ 10¹⁰G. These are too large, though plenty of energy to go around to pump them up if conditions are right [Price-Rosswog, Anderson et.al.,Giacommazo et.al, Liu et.al]

(digress 1) AMR off the way

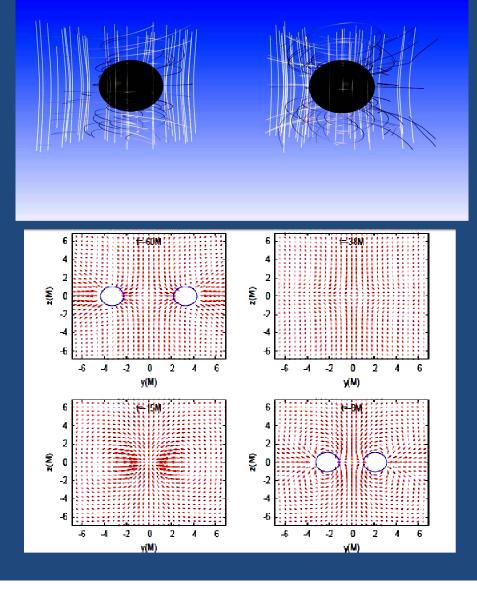
• Tapered approach [LL,Liebling,Reula]

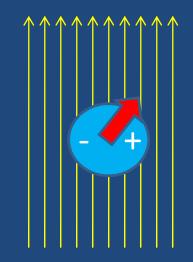






Dynamics... At early times, dynamics 'deduced' from membrane paradigm

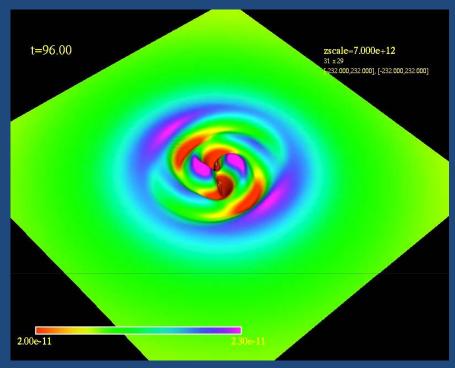




Radiation \rightarrow 2 dipoles in a circular trajectory

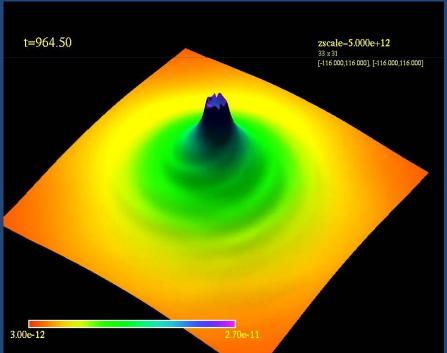
 $E \sim v^3$; $B \sim B_o + v^3$; Flux $\sim v^4$ Quadrupole GW $\sim v^4$

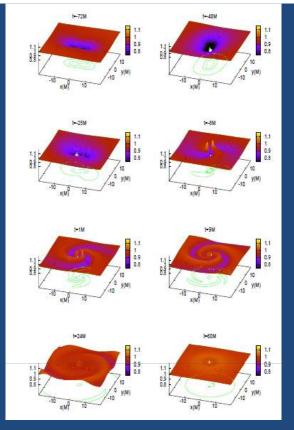
Merger phase

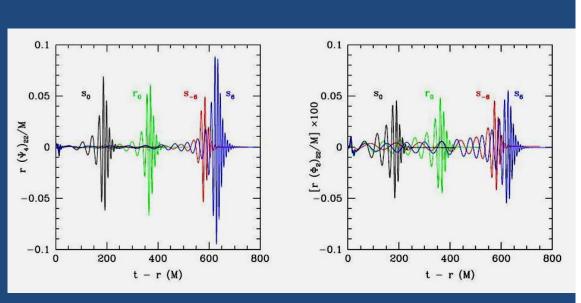


'radial' Poynting

EM Energy density







- Time variability given by orbital motion: 2 dipoles orbiting each other (l=2, m=0 mode). E, B vary as GWs, Poynting 'flux' ½ period. (EM tracers of spacetime?)
- Transition through merger gives rise to I=2,m=2 mode.
 - Induced toroidal electric field
 - $E_T \sim (v_{orbital}/c) B_z \rightarrow Blandford-Znajek analog....?$

Circumbinary disk (pre/prompt)

- Interactions with accretion disk via magnetic fields [GR+ Maxwell]
 - 'Long range' interactions tied by magnetic fields anchored at the disk, reaching the binaries.
 - Binary induces :
 - Strong variability in EM fields, topology change, energy enhancement (reconnection driven effects?)
 - net EM flux can interact with disk , affect accretion rate and emissions by disk. [possible observation needs quiet scenario]
 - Blandford-Znajek (of binary) type mechanism possible though spins required from individual black holes will be high. (spins & B amplitude?)

$$E_{EM} \sim E_{GW} \times 10^{-13} (M/10^8 \text{ Mo})^2 (B/10^4 \text{G})^2$$

QNS:

- Interaction with residual gas yet unexplored
- Individual BZ mechanism?
- Changes due to mass ratios



NS-NS

- Excellent sources of GWs; possible central engine of short GRBs
- Magnetic fields ~ 10¹² G
- GR simulations available, though uncertainties on:
 - eqn of state
 - Role of magnetic fields (err... neutrinos?)
- For grav waves.
 - Early pre-merger stages: PN is good enough
 - Late pre-merger: careful, internal structure may play a role
 - Merger, postmerger: prompt vs. delayed collapse to a BH and other features, we could use to determine eqn of state.
 - Can different effects be disentangled?
- Beyond these, other key qns
 - Does the merger give rise to a BH with sizeable disk?, what is its final spin, magnetic field strength /topology, etc?
 - How long does the hypermassive star exists before collapsing?
 - All these connect directly with short GRBs models

NS-NS.

- Einstein equations
 - Generalized Harmonic formulation:
 - Constraints :
 - Einstein eqns:

$$C_a = \Gamma_a + H_a$$

$$R_{ab} = \nabla_{(a}C_{b)} + {}^{TR}T_{ab} + \kappa \{2n_{(a}C_{b)} - g_{ab}n^cC_c\}$$

- GRHydro:
 - Eqns determined by:

$$\nabla_a T^{ab} = 0 \quad ; \quad \nabla_a (\rho u^a) = 0$$

$$T_{ab} = (\rho_0 (1 + \varepsilon) + P) u_a u_b + P g_{ab} + F_a^c F_{bc} - \frac{1}{4} F^{cd} F_{cd}$$

 $\nabla^{a} \nabla_{a} x^{u} = H^{u}$

- Expressed in terms of conservative variables, (use of HRSC)
- Eqn of state: $P = (\Gamma 1)\rho_o \varepsilon$ (though, $P = k\rho_0^{\Gamma}$ for ID)
- ideal MHD limit
 - Constraint (s)?
 - How to deal with it? (eqns weakly hyperbolic if non zero, or something else done)
 - How to not break it when using AMR? \rightarrow tapered method
 - How to ensure boundaries won't get in the way? \rightarrow constraint pres. bound cond.
 - How to make all transparent for global treatment of the problem?

(digress 2) MHD eqns. Simple example

α = 0 : conservative,
but weakly hyperbolic
α = 1 : non-conservative,
but hyperbolic!
c₁ ; s drivers of constraint

$$\begin{split} \partial_t \rho &= -\nabla_i (\rho v^i) \\ \rho \partial_t v^i &= -\rho v^j \nabla_j v^i - \nabla^i p - B_k (\nabla^i B^k - \nabla^k B^i) - \alpha B^i \nabla_k B^k \\ \partial_t B^i &= -\nabla_j (u^j B^i - u^i B^j) - \alpha u^i \nabla_j B^j - c_l \nabla^i \phi \\ \partial_t e &= -\nabla_i ((e + p + \frac{1}{2} B^2) v^i - B^i v \cdot B) - \alpha v^i B_i \nabla_k B^k - c_l B^k \nabla_k \phi \\ \partial_t \phi &= -\alpha u^j \nabla_j \phi - c_l \nabla_j B^j - s \phi \end{split}$$
 where

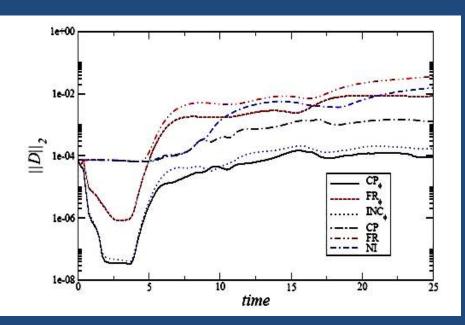
$$p := (\gamma - 1)(e - \frac{1}{2}\rho v^2 - \frac{1}{2}B^2) \qquad c_s^2 := \gamma p/\rho$$

• Strongly/symmetric hyperbolic \rightarrow complete set of eigenvectors

Constraint Preserving Bound. Cond.
→induce eqn for boundary values

$$(\nabla B)_{t} = 0 \rightarrow (\partial_n B)_{t} = known + bdry data$$

[Cecere,LL,Reula 07]



[Cecere,LL,Reula]

Constraint control

Boundary conditions

Flexibility of picking algorithms (eg. not tied to constraint transport)

GRMHD. Star + by-hand violation

[Liebling,LL,Neilsen,Palenzuela]

(digress 3) Last rabbit out of a hat

'Extended' ideal MHD equations

- Terms added for :
 - Divergence cleaning (c_r , c_h)

• Ensure strong hyperbolicity eve if no div cleaning used.

 if no added field, 'eight' wave formulation. (but has sources with derivatives) $\partial_t \tilde{D} + \partial_i \left| \alpha \tilde{D} \left(v^i - \frac{\beta^i}{\alpha} \right) \right| = 0,$ $\partial_t \tilde{S}_j + \partial_i \left[\alpha \left(\tilde{S}_j \left(v^i - \frac{\beta^i}{\alpha} \right) + \sqrt{h} P h^i_j \right) \right]$ $= \alpha^{3} \Gamma^{i}{}_{jk} \left(\tilde{S}_{i} v^{k} + \sqrt{h} P h_{i}{}^{k} \right) + \tilde{S}_{a} \partial_{j} \beta^{a}$ $-\partial_i \alpha (\tilde{\tau} + D)$ $-\zeta \alpha (\tilde{B}_i W^{-2} + v_i v_j \tilde{B}^j) \partial_k \tilde{B}^k$, $\partial_t \tilde{\tau} + \partial_i \left[\alpha \left(\tilde{S}^i - \frac{\beta^i}{\alpha} \tilde{\tau} - v^i \tilde{D} \right) \right]$ $= \alpha \left[K_{ij} \tilde{S}^i v^j + \sqrt{h} K P - \frac{1}{\alpha} \tilde{S}^a \partial_a \alpha \right],$ $-\zeta \alpha v_i \tilde{B}^j \partial_k \tilde{B}^k$ $\partial_t \tilde{B}^b + \partial_i \left[\tilde{B}^b \left(v^i - \frac{\beta^i}{\alpha} \right) - \tilde{B}^i \left(v^b - \frac{\beta^b}{\alpha} \right) \right]$ $= -h^{bj}\partial_{j}\Psi - \alpha\sqrt{h}h^{ij}\partial_{j}\Psi - \zeta\alpha v^{i}\partial_{j}\tilde{B}^{j}$ $\partial_t \Psi = -c_r \alpha \Psi - c_h \frac{\alpha}{\sqrt{h}} \partial_i \tilde{B}^i + (\beta^i - \alpha v^i) \partial_i \Psi$

Initial configuration. (Not 'too physical'...)

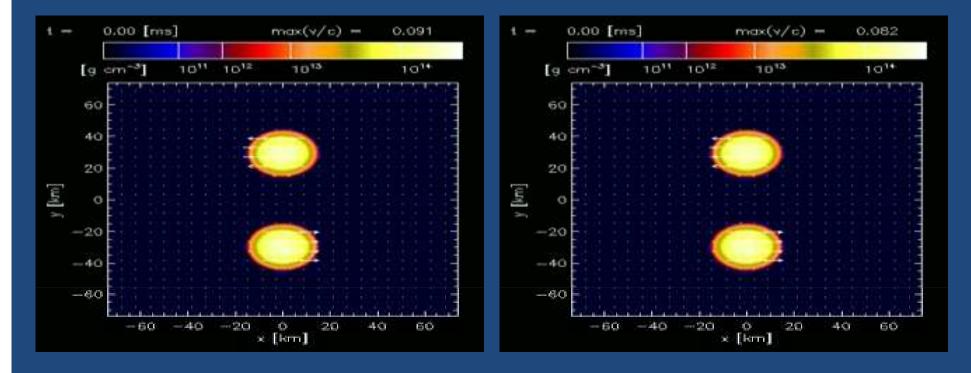
- Equal non-rotating polytropes to represent the stars (Γ =2). (R_s=16.26km, M~0.9 M_o)
- Poloidal seed magnetic fields, antialigned with orbital angular momentum in each star. B ~ 9 10¹⁵ G

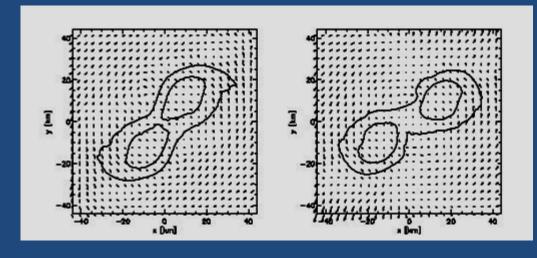


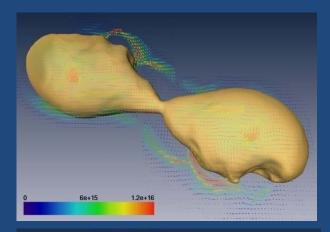
- Initial separation ~ 4 R_s
- Grid : $[-100R_s, 100R_s]$; up to 7 levels of refinement Δ_{min} =0.46km. Gravitational waves extracted well within the wave zone.



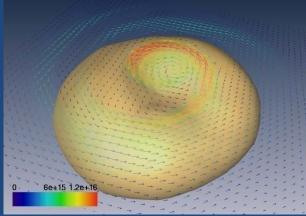
[Anderson, Hirschmann, LL, Motl, Neilsen, Palenzuela, Tohline 08]







 $\rho = 10^{14} g/cm^3$



ρ=8 10¹³g/cm³

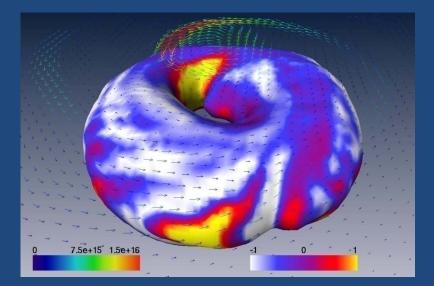
 K-H instability, shear energy into magnetic fields which grow ~ 1 order of magnitude. Growth saturates ~ Alfven time-scale

 Hypermassive star differentially rotating at rates different from the non-magnetized case

0 Be+15 1.2e+16

 $\rho = 10^{14} \text{g/cm}^3$

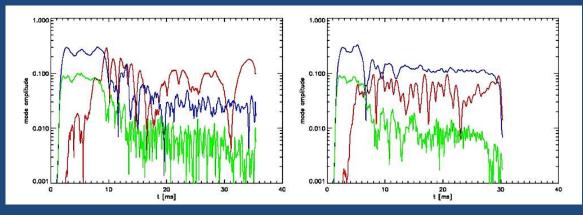
Less obvious effects



Magnetic Buoyancy



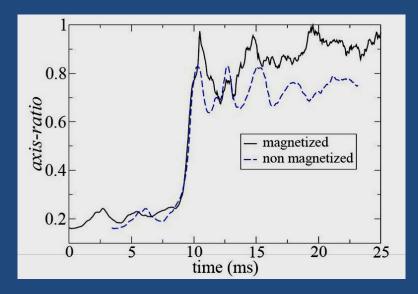
B ~ 10¹⁷G

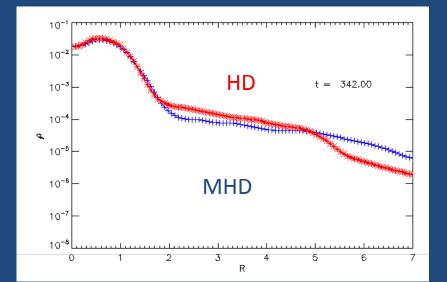


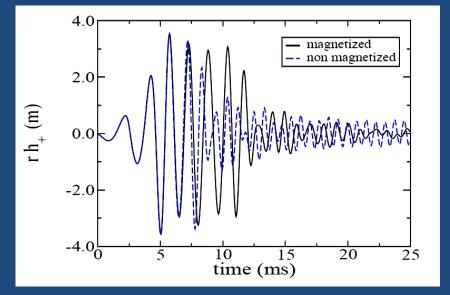
I = 2 decreasing consistently with gravitational wave emission rateI = 4 'Cartesian grid' induced remains under control throughout

I = 1 mode growing, consistent with Tayler's instability (magnetic field)

'bulk' dynamics & consequences







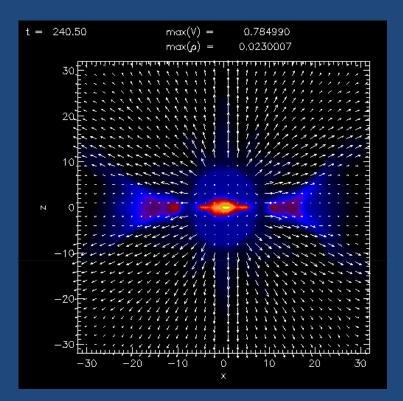
Amplitudes ~ Consistent with noise level to 2-4Mpc in current LIGO

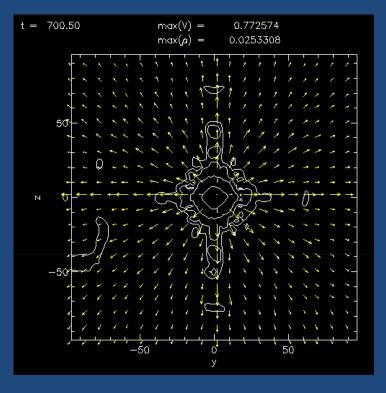
Matches:

- 0.98 to t=7.5ms
- 0.63 to t= 12.5ms
- 0.59 to t=25ms

[related to Price-Rosswog 06]

Other hints...





• Could these induce asymmetric bubbles? [Bucciantini et. al.]

• how about long-term behavior of the disk? Other models will need to take over [Metzger et.al.]

•No right to expect this would work! Not the right approx to study the different regimes

as opposed to BH-BH case, too early to hand it over....

BH-NS

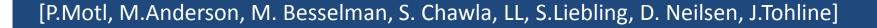
- GR studies ongoing (NS: polytrope or via Shen eos)
 - Sizeable disks if sufficiently high spins (>~ 0.5)
 - 'competition' with mass ratio
 - Magnetic field effects?

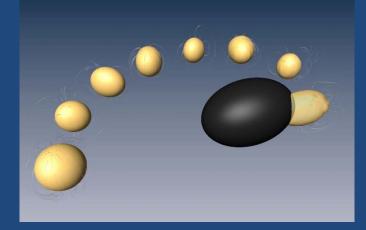
Neutron Star: Irrotational, $\Gamma = 2$ R = 15 [km] M = 1.4 M_{solar}.

Initial dipole field of strength 10¹² [Gauss]

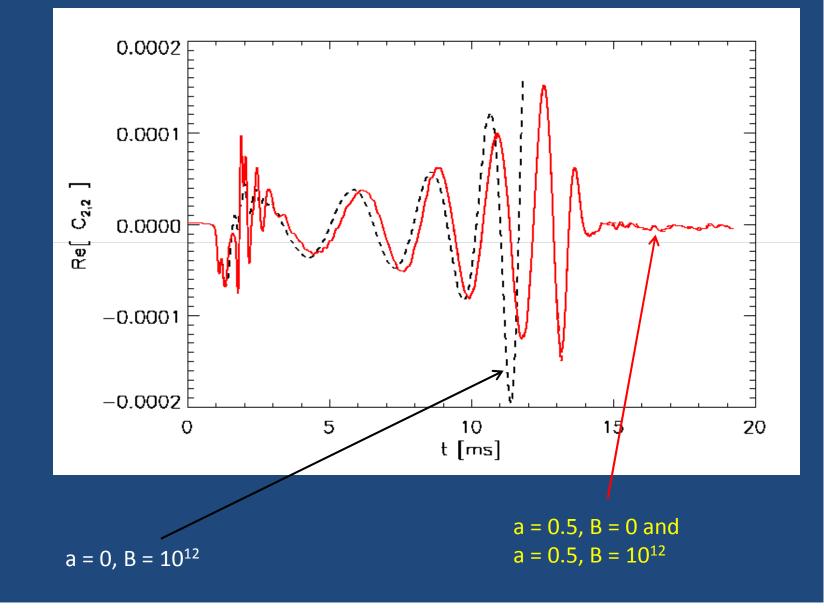
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Black Hole: M = 7 M_{solar}; a/M = 0, 0.5
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Initial separation of 100 [km] Grid extends to ± 443 [km] Peak resolution of 0.73 [km] or 40 points across initial neutron star

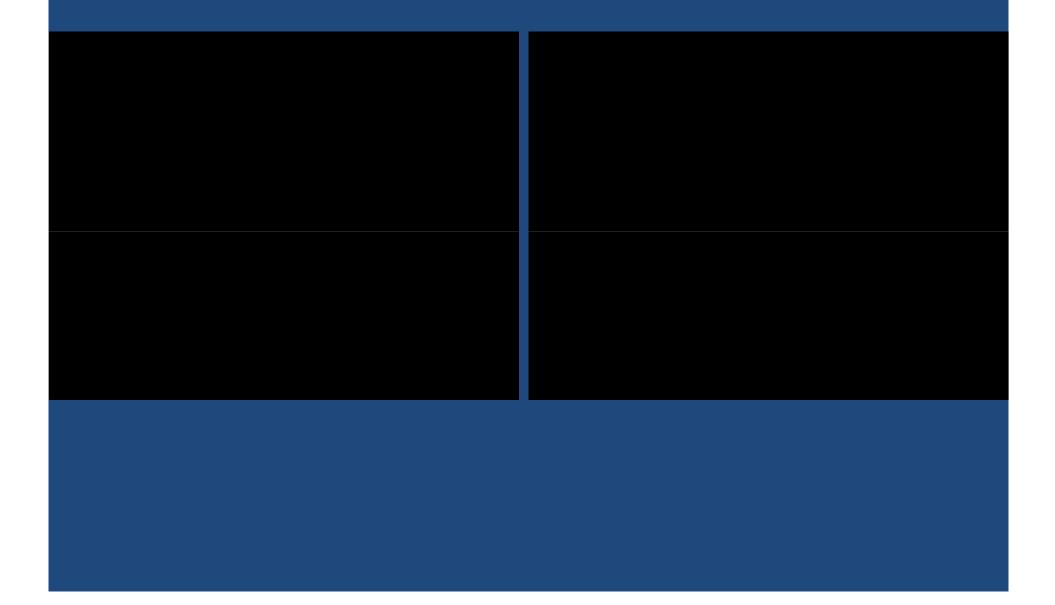




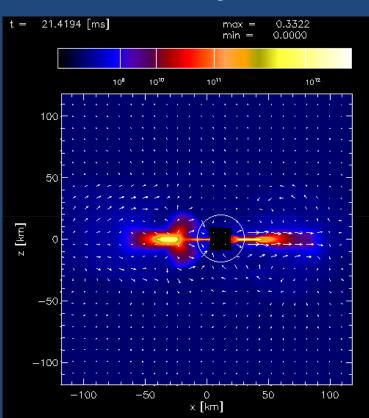
Grav. Waves...



Different cases...

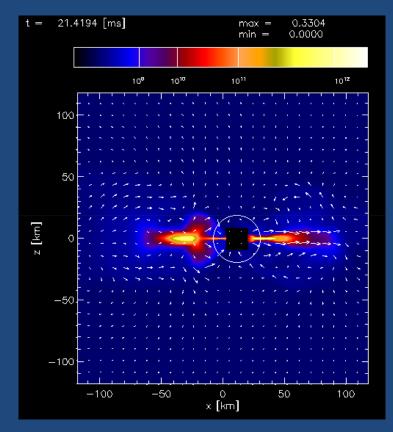


Disk and structure (a/M=0.5)



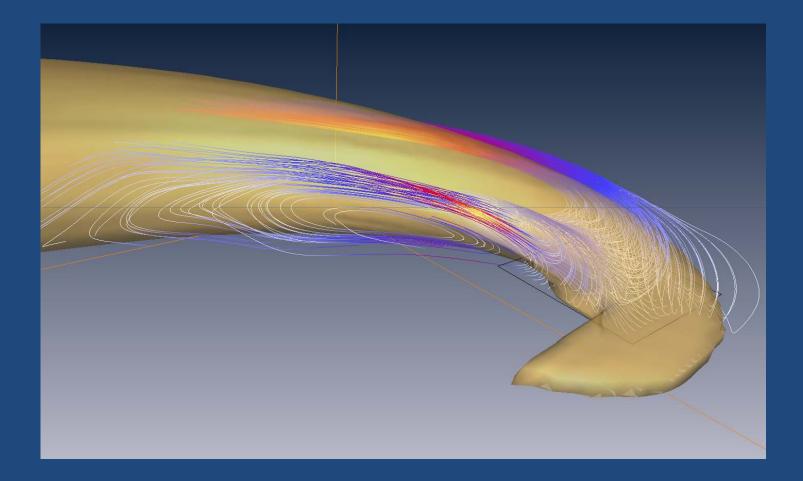
Unmagnetized

$B = 10^{12}$ initially



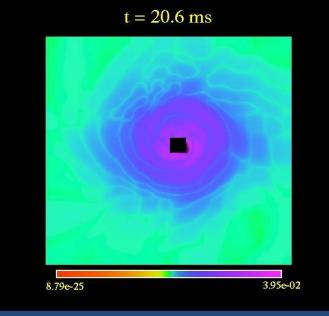
1.6%

Field lines & behavior



Handing it over...



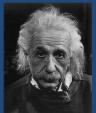


[t=15.6 ms] → Going toroidal... [t=20ms] B ~ $10^{12} \rightarrow 10^{13}$ G

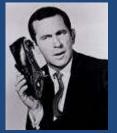
--Long times.... Radiation transport effects to be accounted for yet --connection with fall back model for SGRB's tails [Rosswog,Lee&Ramirez Ruiz]

Final words

- Interesting problems can be studied with different approx and connect with other studies
- Full problem requires ability to study several regions at the same time that will change dynamically.
 - Building up a new approach. GR + Maxwell + Fluid eqns with suitable model for currents, lying half way in between:



- Novel, unorthodox approach through the path less taken
- if brute force doesn't work... use more of it!



Talk by Palenzuela on Saturday.