

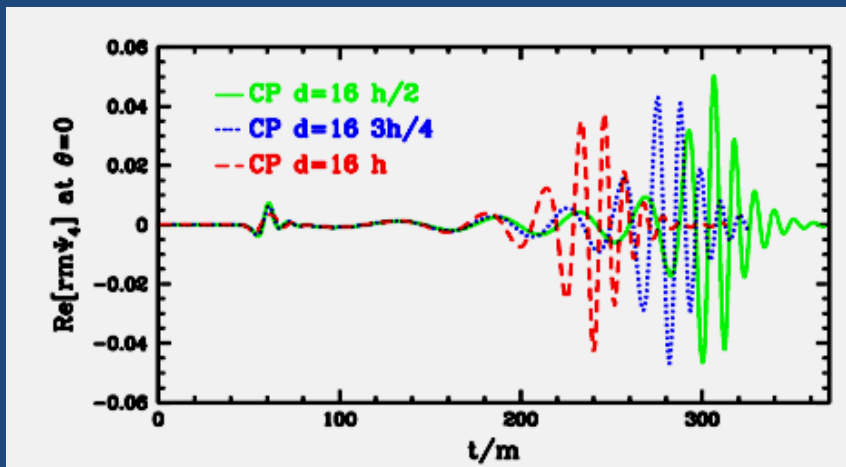
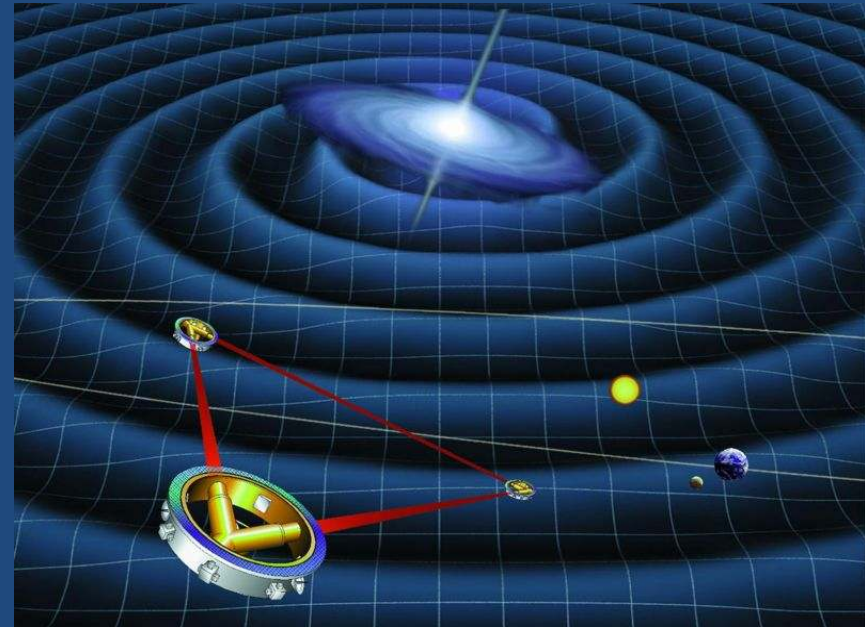
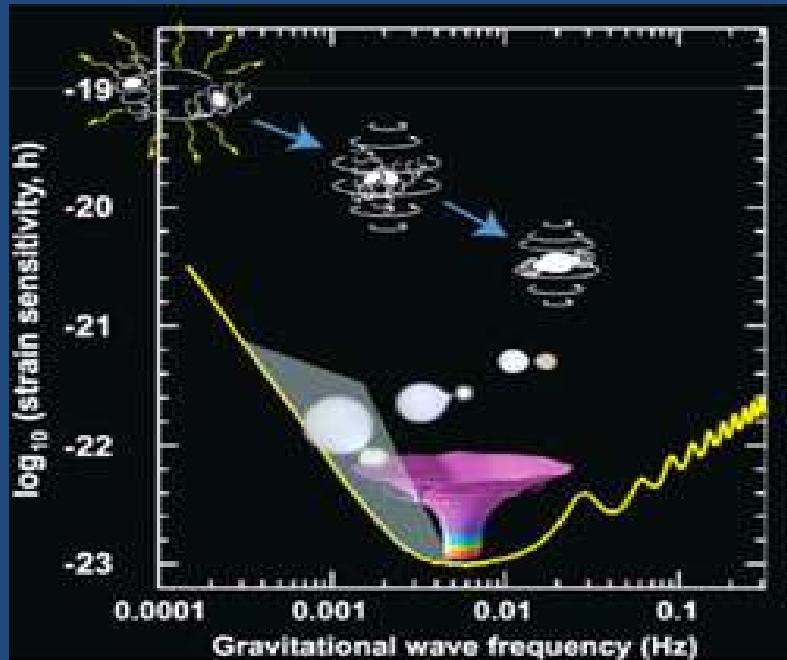
Electromagnetic counterparts to loud gravitational wave events

L. Lehner

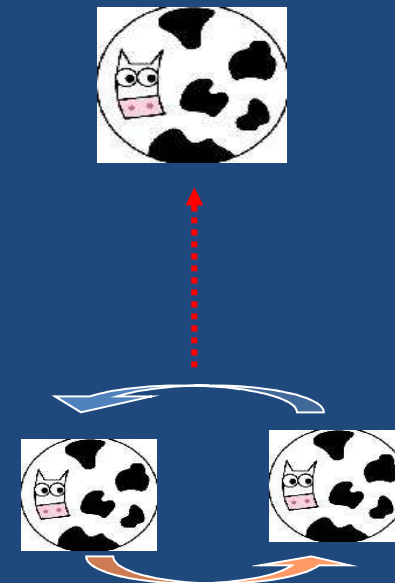
(Perimeter Institute/Univ. of Guelph/CIFAR)



Driving goal I



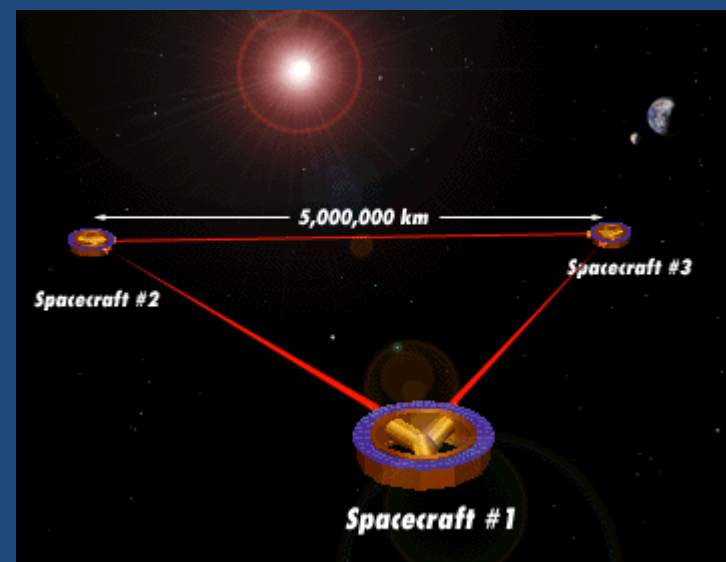
[Pretorius 05,]



Not so fast?

LISA: superb signal to noise ratio
[Phinney's, Baker's talks]

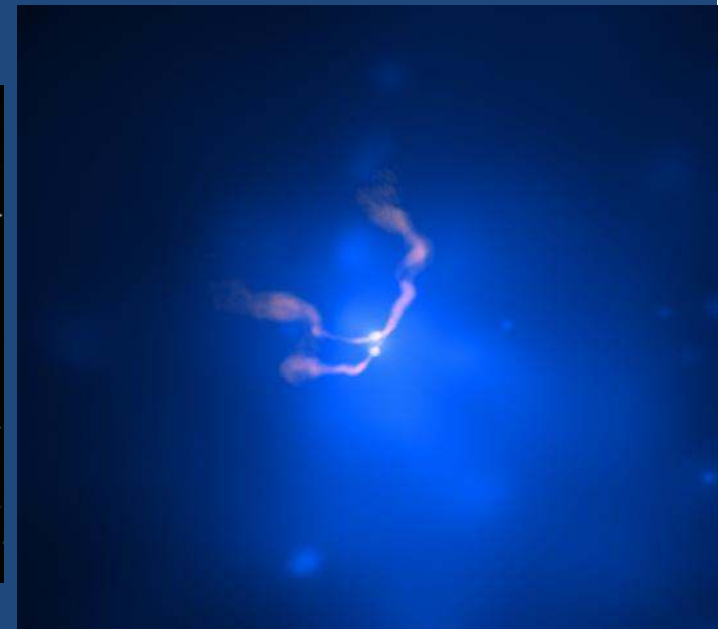
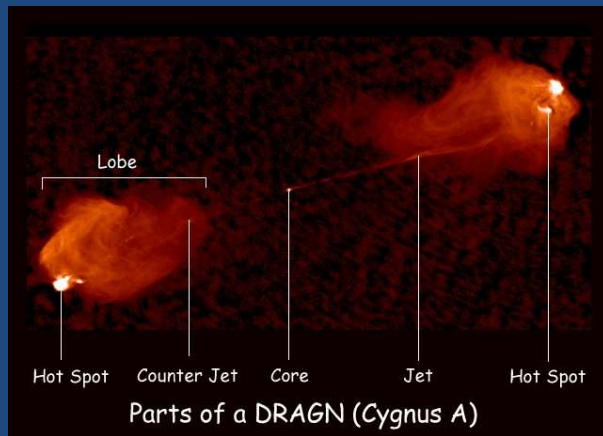
- waves will be ``seen'' directly and to very large redshifts ($z \sim 5-10$...)
- Potential to do accurate astrophysics
- Potential to probe cosmology
- Potential to probe higher dimensional scenarios
- However:
 - localization to \sim square degrees [Cutler's talk; Holtz-Hughes]
 - distance obtained is redshift dependent



An electromagnetic counterpart resolve these issues

Nature cooperates...

- Understand both gravitational and electromagnetic wave emissions from key systems
 - Binary black holes interacting with surrounding media



10^{51} Ergs routinely inferred... even some 10^{61} ergs ?!

Studying relevant systems (BBH)

- Deal with spacetime curvature
 - Einstein equations. That's the 'solved' part! (ie... if you 'think' about it.. NR can give the answer)
- Black holes... are not really quite in vacuum...must deal with fields describing gas and electromagnetic fields
 - Poorly understood systems [we don't control the experiment]
 - Matter, what matter ?
 - Electromagnetic fields?
 - Emission process?



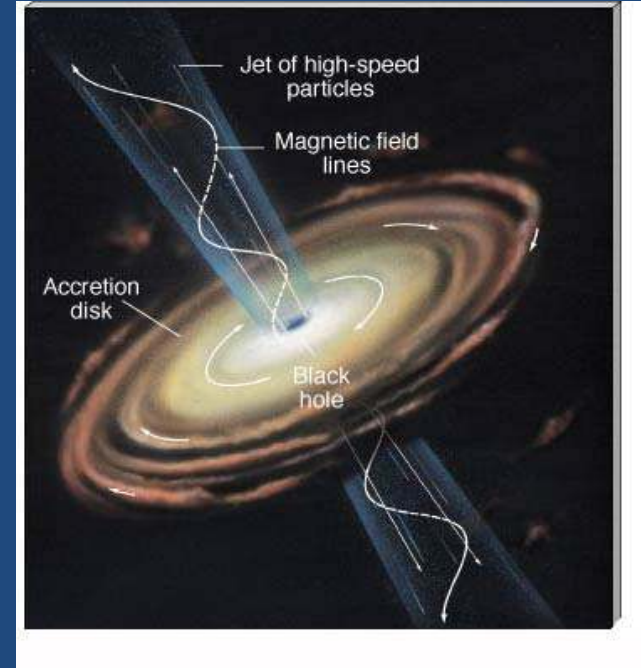
Two fronts.

(circumbinary picture [complementary to Bogdanovic])

- Pre/prompt/post - merger emissions?
 - (pre/prompt) Binary black holes as EM field stirrers
 - (post) Binary black holes as bullies for matter

Merger of galaxies

- observations indicate the presence of **supermassive BHs in the center of galaxies**, surrounded by gas and an accretion disk
- these galaxies have undergone mergers → binary black hole merger
- further, AGNs → **BHs are surrounded by a disc of matter likely magnetized.**



Binary black holes and emissions

- Different possible options.
 - Postmerger events from circumbinary disks around BHs



[Milosavljevic-Phinney;
Lipai-Loeb;
Lipai et.al,
Bonning et.al;
Bode et.al;
O'Neil et. al;
Megevand et.al;
Corrales et.al, etc.]

- Pre/merger events from gas/plasmas in between BHs /
torques on disk



[Armitage et.al;
MacFadyen et.al.;
Dotti et.al;
Chang. et.al.;
Palenzuela et.al.;
Bode et.al...]

After merger consequences

- Circumbinary disk ‘knows’ a merger takes places ‘after the fact’
 - $\sim 5\%$ energy radiated, most during last orbit: gravitational potential weakens ‘suddenly’
 - Recoil in a given direction
- in both cases, the disk needs to readjust

Approach: GR+Hydrodynamics

- Einstein equations

- Generalized Harmonic formulation:

$$\nabla^a \nabla_a x^u = H^u$$

- Constraints : $C_a = \Gamma_a + H_a$

- Einstein eqns: $R_{ab} = \nabla_{(a} C_{b)} + {}^{TR}T_{ab} + \kappa \{ 2n_{(a} C_{b)} - g_{ab} n^c C_c \}$

- Expressed in terms of g_{ab} ; H_a from stationary solution.

- *Cowling approximation is enough.*

- GRHydro:

- Eqns determined by:

$$\nabla_a T^{ab} = 0 \quad ; \quad \nabla_a (\rho u^a) = 0$$

$$T_{ab} = (\rho_0(1 + \varepsilon) + P)u_a u_b + P g_{ab} + F_a^c F_{bc} - \frac{1}{4} F^{cd} F_{cd} g_{ab}$$

- Expressed in terms of *conservative variables*, (use of HRSC)

- Eqn of state:

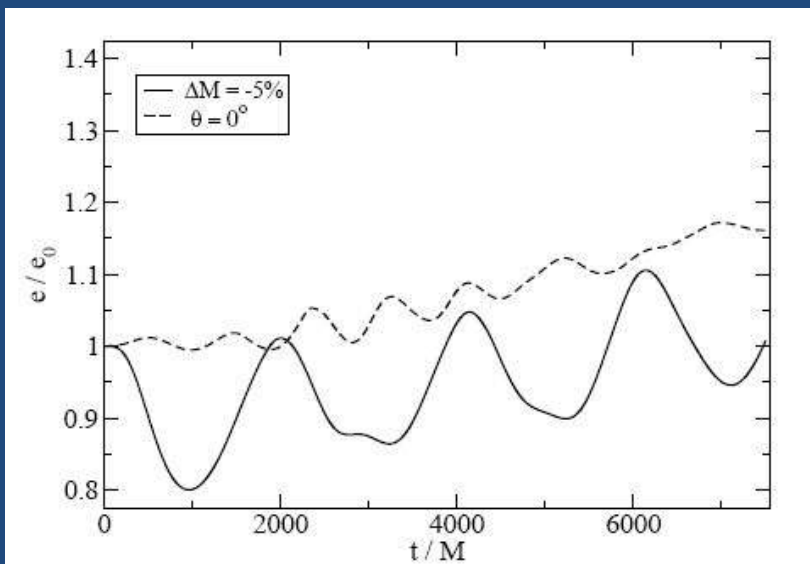
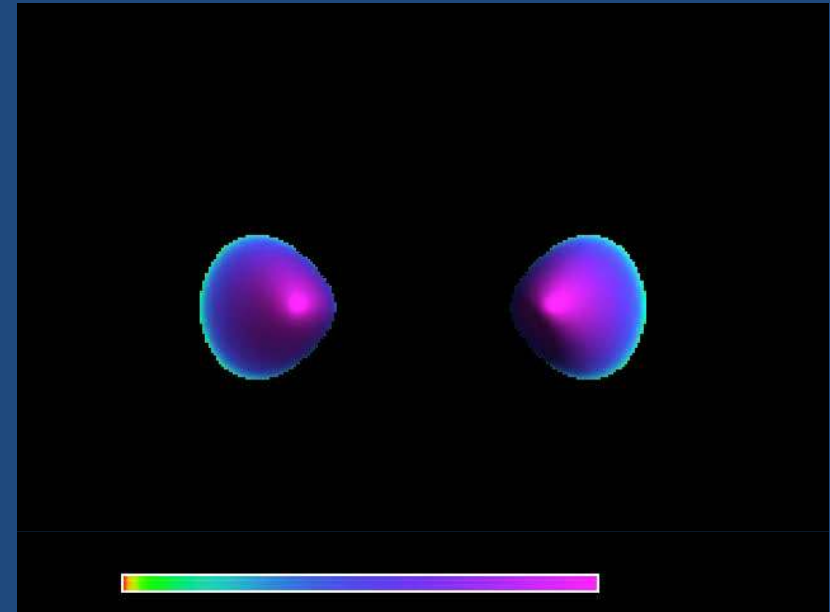
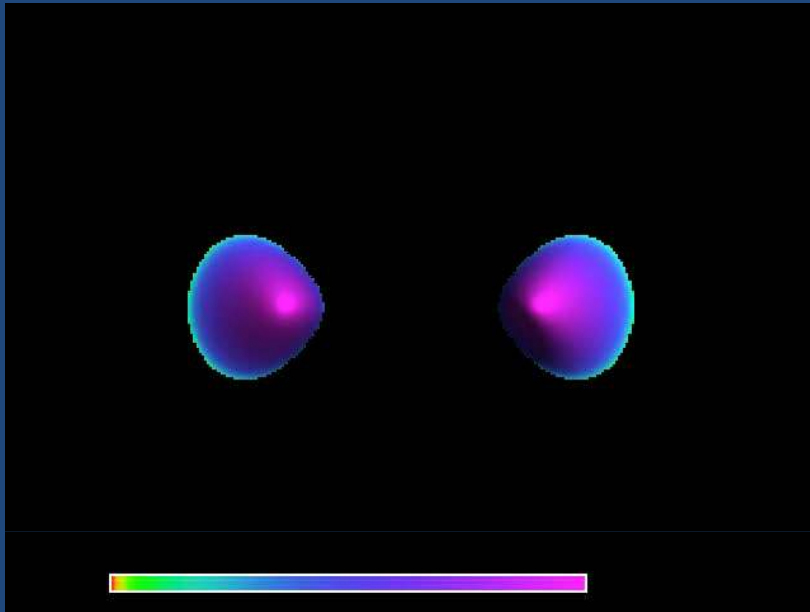
$$P = (\Gamma - 1)\rho_0 \varepsilon \quad (\text{though, } P = k\rho_0^\Gamma \text{ for ID})$$

- Present case $\Gamma=5/3$, ideal MHD eqns. Results for $B=0$.

‘kicked black holes’ retaliate

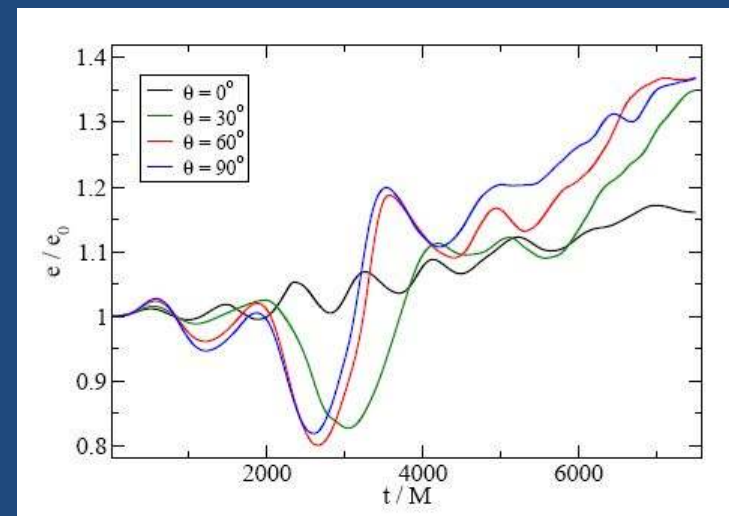
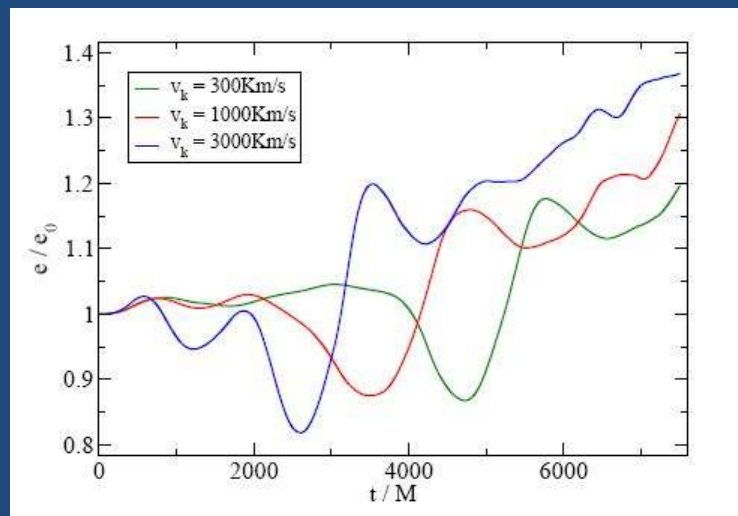
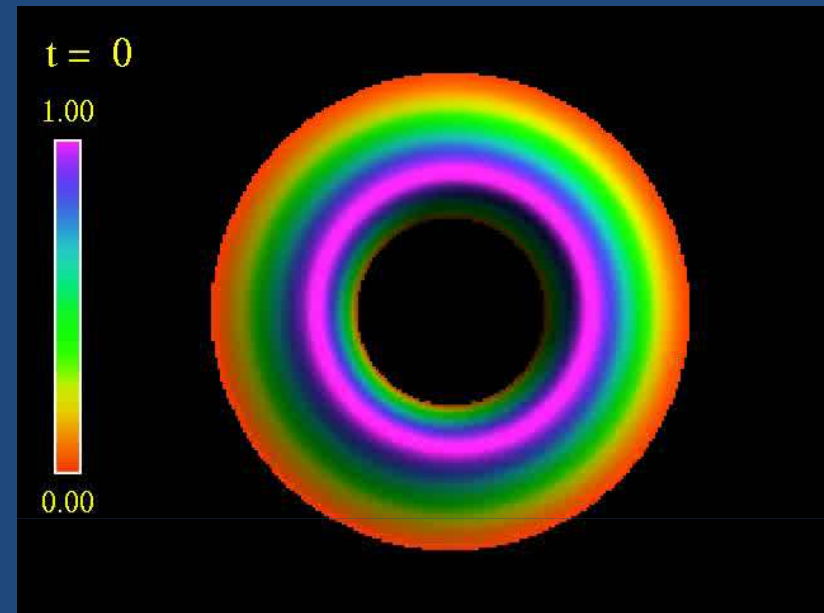
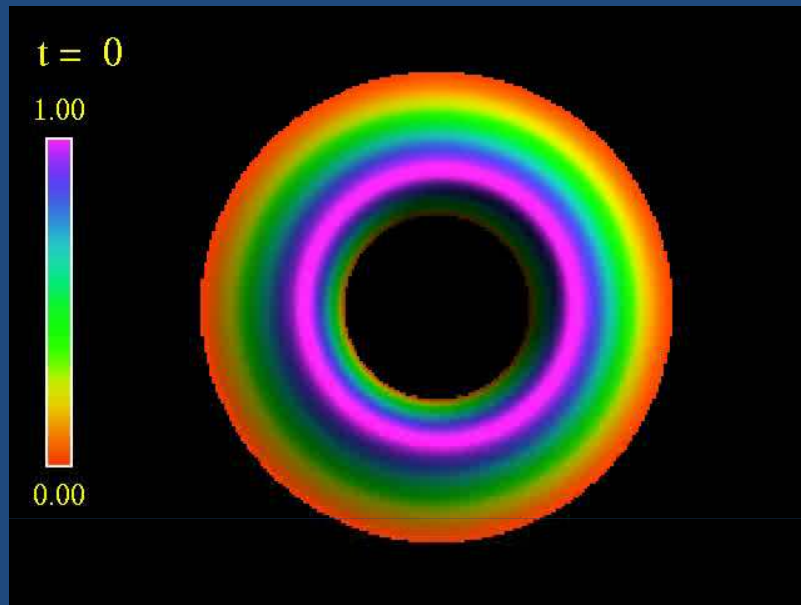
- Both mass reduction and recoil speed have an impact on the disk’s behavior. Relevant quantities: v_{kick} ; v_{sound} ; v_{orbital}
- If shocks develop → shock energy onto the disk → can induce EM signals
 - *from there on... take your pick...*
 - Lipai et. al. : prompt and in the UV
 - Bonning et. al. : delayed and in soft Xrays
 - Phinney et. al. : not kicks but mass reduction, significant output
 - O’Neil et. al. : not kicks but mass reduction, lowering of luminosity

Symmetry preserving cases



- Internal energy reduction (same with pressure, temperature).
- Possible reduction in disk luminosity initially, but oscillates.
- Time variability governed by the disk's period

Symmetry breaking cases



- In all cases, significant heating of disk is induced
 - Robust behavior
 - Time of ‘swing’ appearance
 - $T \sim 5200 - 912 \ln ((v_k)_{\rightarrow} / 300 \text{ km/s})$
 - Strong (short) variability would certainly impart its characteristics to EM production process
 - Caveats:
 - Long term behavior, influenced by Papalaziou-Pringle ‘instability’
 - $T \sim T_{\text{ref}} - (2/w) \ln ((v_k)_{\rightarrow} / v_{\text{ref}})$, $w \sim 0.45$ (too high!)
 - Ongoing work on thin disk to remove this issue.
 - Role of Magnetic fields might be important.
 - Role of cooling important!
 - Particular emission process?

Taking images..

Radiation transfer eqn:

$$\frac{d I_\nu}{d\lambda} = -p^a u_a (\eta_0 - \chi_0 I_\nu)$$

Options:

Brehmstrahlung-blackbody model

η brehmstrahlung emmisivity

χ modified Krammer's opacity law

Thermal model

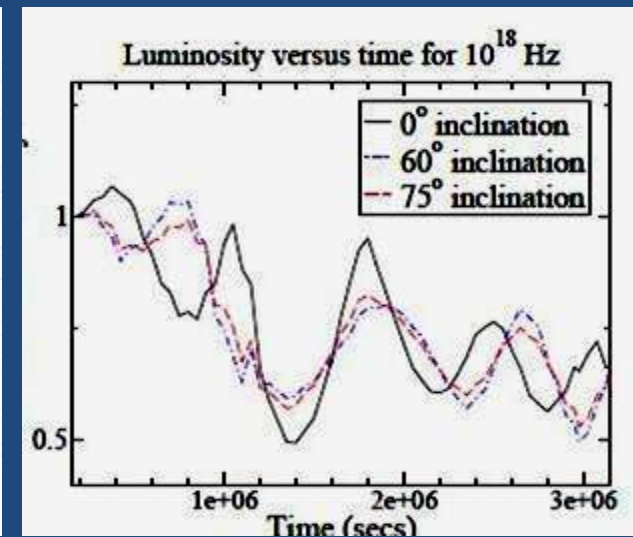
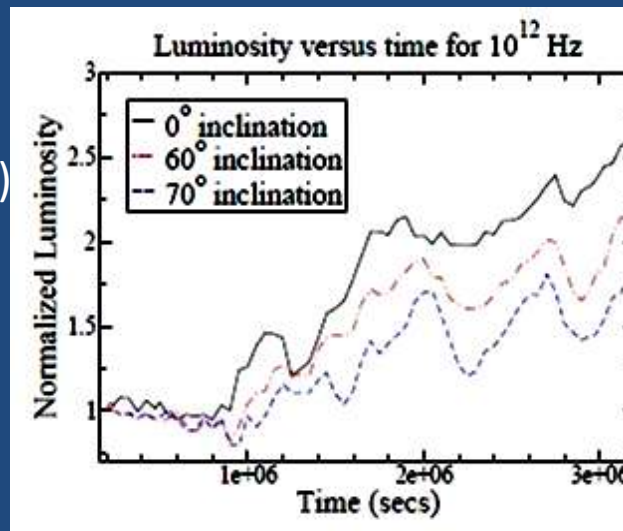
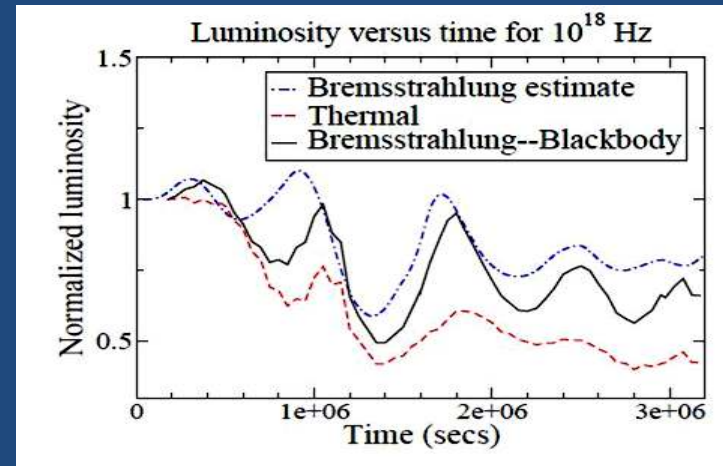
η Kirchoff's law $\eta=\chi$ (Planck law)

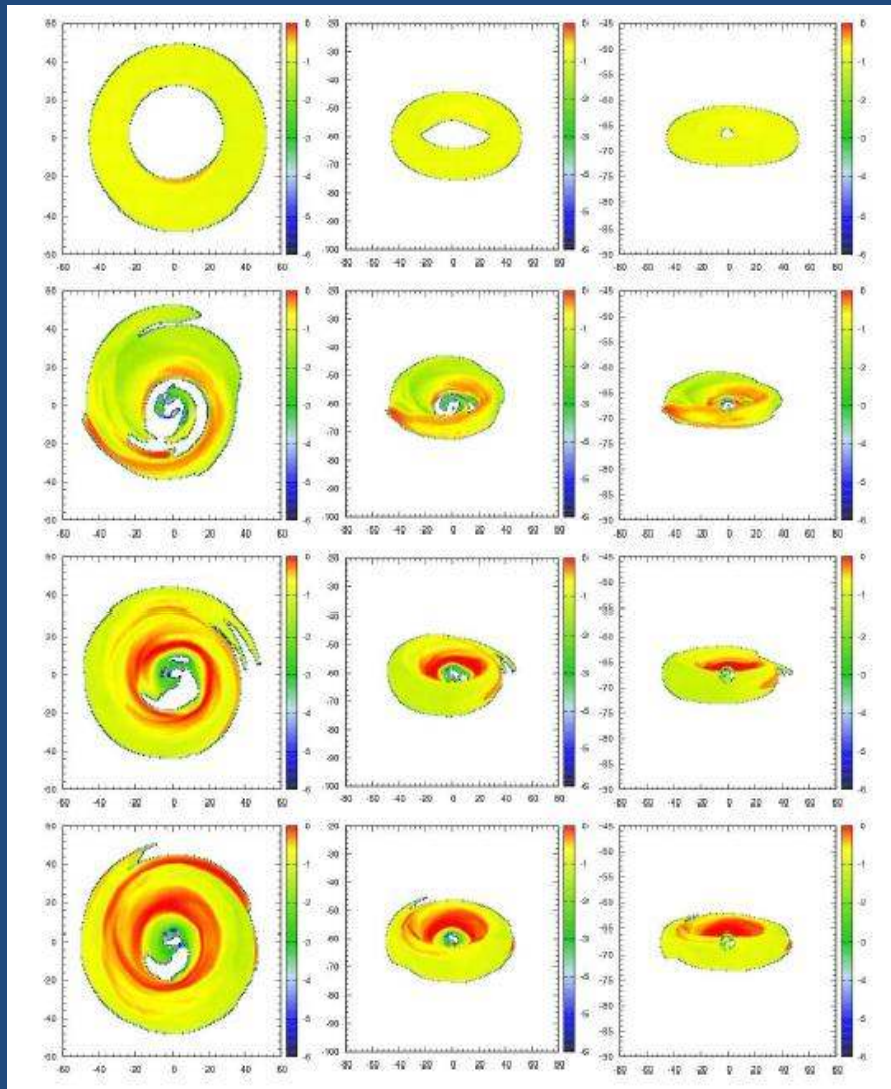
χ Krammer's opacity law

Brehmstrahlung vanilla

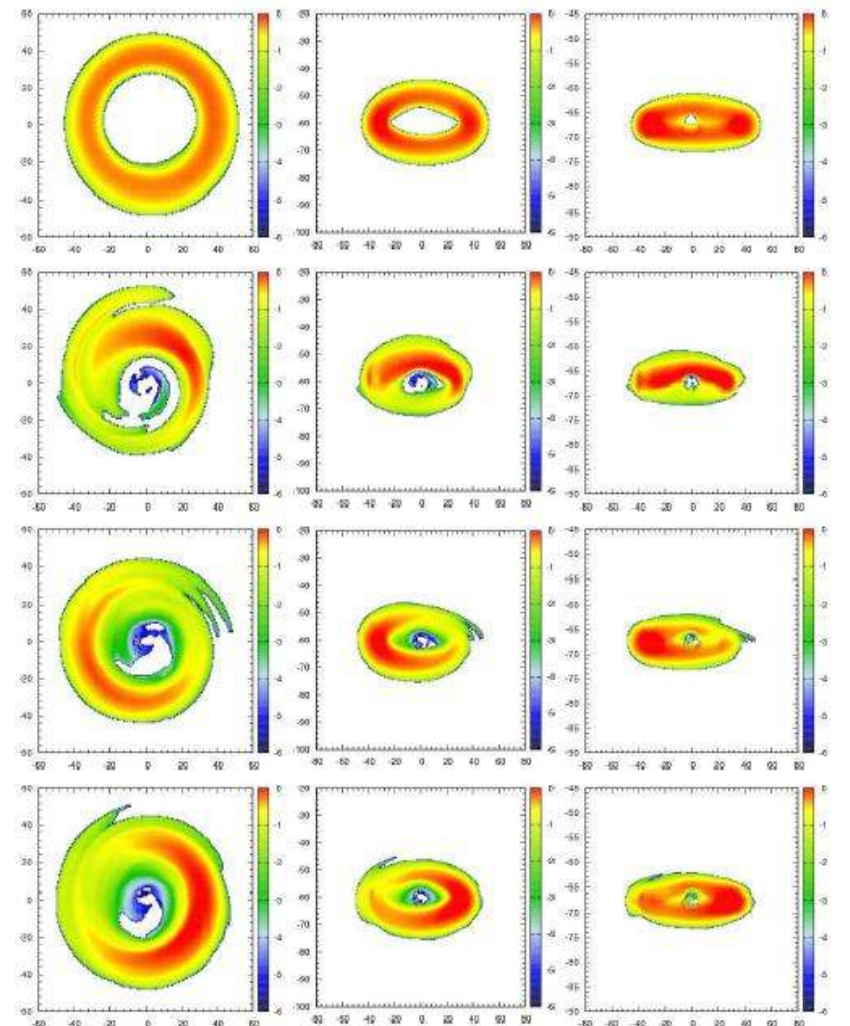
$$\int \rho^2 T dV$$

[a la, Schnittman; Noble]





Infrared



Xrays

[Anderson et.al, 2009]

Conclusion (I)

- BH recoil can produce observable consequences by affecting the disk.
- However... it might be too late, need to roll a dice with 10^5 faces!

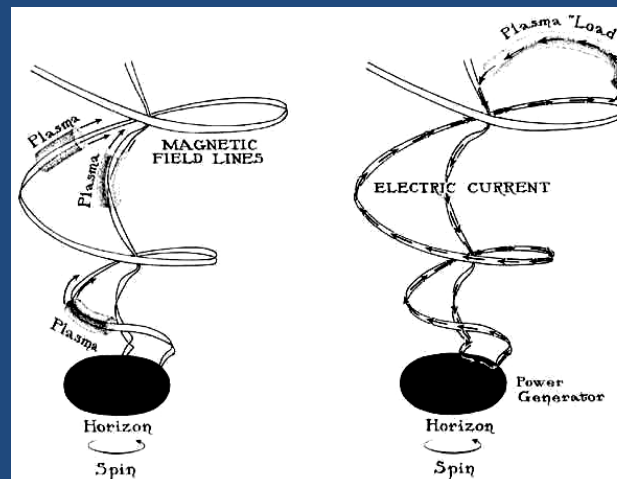
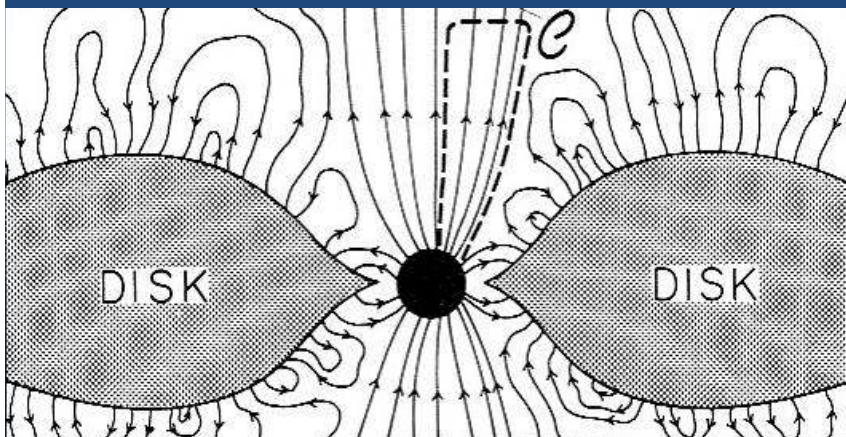
Binary black holes as blenders.

A new spin on an old story (though without spin)

How does the curvature/dynamics influence EM fields?



- Blandford-Znajek. “Penrose” process for Kerr bh’s surrounded by magnetic fields (anchored by the disk)
- Stray charges accelerate \rightarrow pair production cascade. BH becomes surrounded by a tenuous conducting plasma with little inertia



[Goldreich-Julian,
Blandford-Znajek]

Approach: Force-free electrodynamics

$$\nabla_a T^{ab} = 0 \quad \rightarrow \quad \nabla_a T^{ab}_{(\text{fluid})} = -\nabla_a T^{ab}_{(\text{em})} = -F^{ab} J_a$$

$$\text{if } \rho, P \ll B^2 \quad \text{then} \quad \nabla_a T^{ab}_{(\text{fluid})} \ll F^{ab} J_a \approx 0$$

$$\mathbf{E} \cdot \mathbf{J} = 0, \quad \mathbf{q} \mathbf{E} + \mathbf{J} \times \mathbf{B} = 0 \rightarrow \mathbf{E} \cdot \mathbf{B} = 0$$

Stationary spacetime:
(Gammie, McKinney 04)

$$E \bullet B = 0 \rightarrow F_{ab} * F^{ab} = 0 \rightarrow A_{\phi,\theta} A_{t,r} - A_{t,\phi} A_{\phi,r} = 0$$

so

$$\Omega_F \equiv - \frac{A_{t,r}}{A_{\phi,r}} = - \frac{A_{t,\theta}}{A_{\phi,\theta}} = \frac{F_{tr}}{F_{r\phi}} = \frac{F_{t\theta}}{F_{\theta\phi}}$$

$$E_{,t} = 2\pi \int_0^\pi \sqrt{-g} F_E d\theta \quad (\text{with } F_E = -T_t^r)$$

$$\rightarrow F_E = 2(B^r)^2 \Omega_F r \left(\frac{a}{2Mr} - \Omega_F \right) \sin^2(\theta) - B^r B^\theta \Delta \sin^2(\theta)$$

thus,

$$F_E \Big|_{r=r_H} = 2(B^r)^2 \Omega_F r_H (\Omega_H - \Omega_F) \sin^2(\theta)$$

$$\Rightarrow \text{for } 0 < \Omega_F < \Omega_H \quad \text{and} \quad B^r \neq 0 \quad \text{energy out of horizon}$$

Plasma is crucial for
this to happen

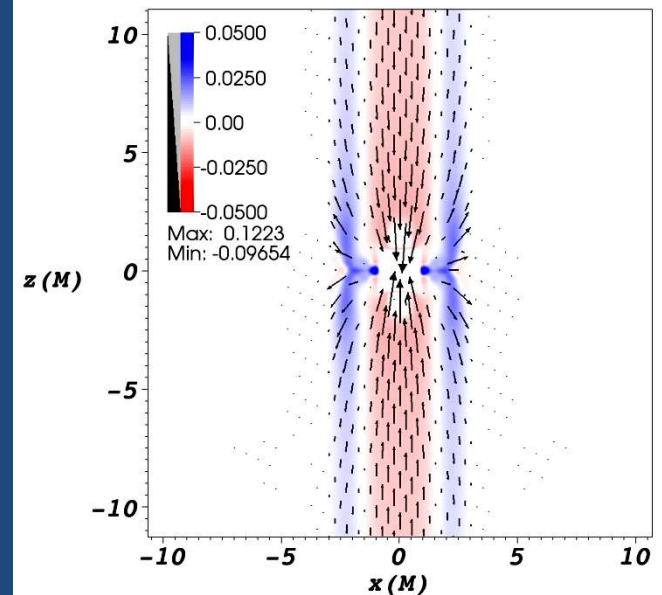
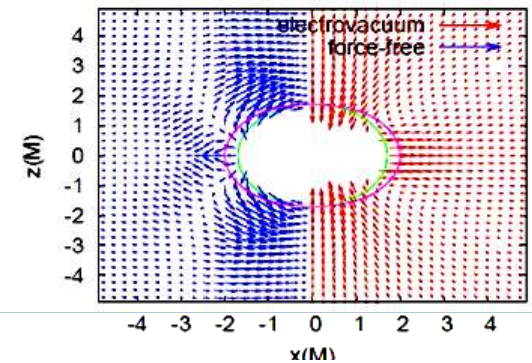
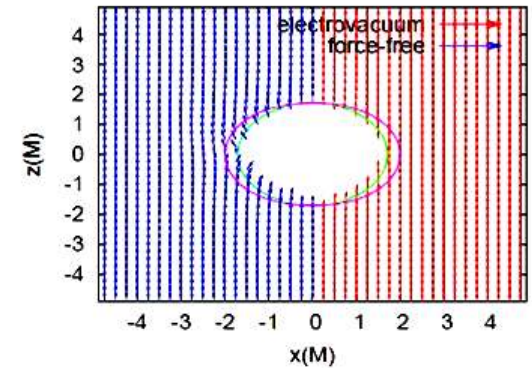
Examples...

- Kerr in vacuum and FF immersed in uniform field
- In vacuum \rightarrow no radiation
- With plasma \rightarrow currents on the horizon 'complete the circuit'

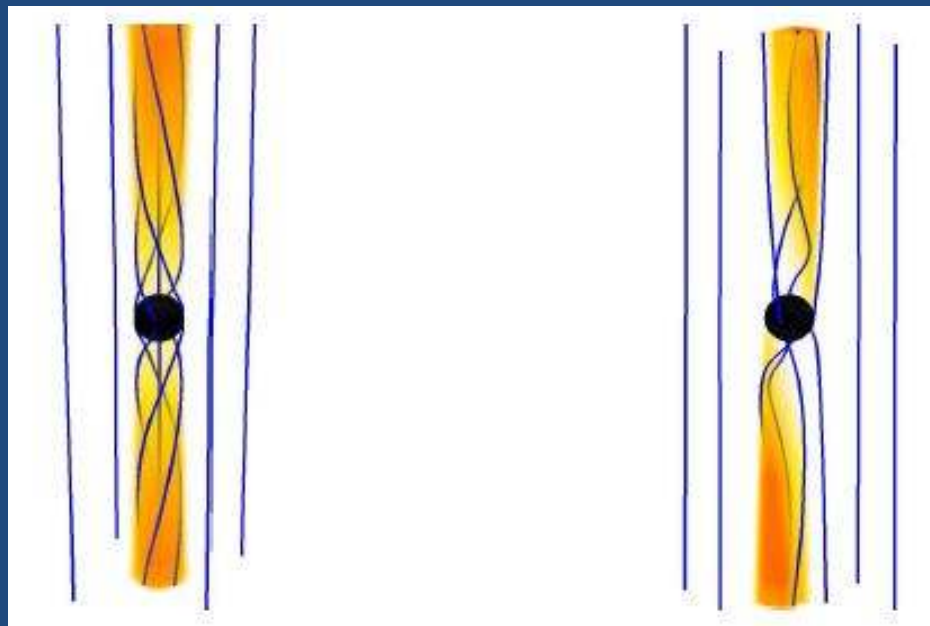
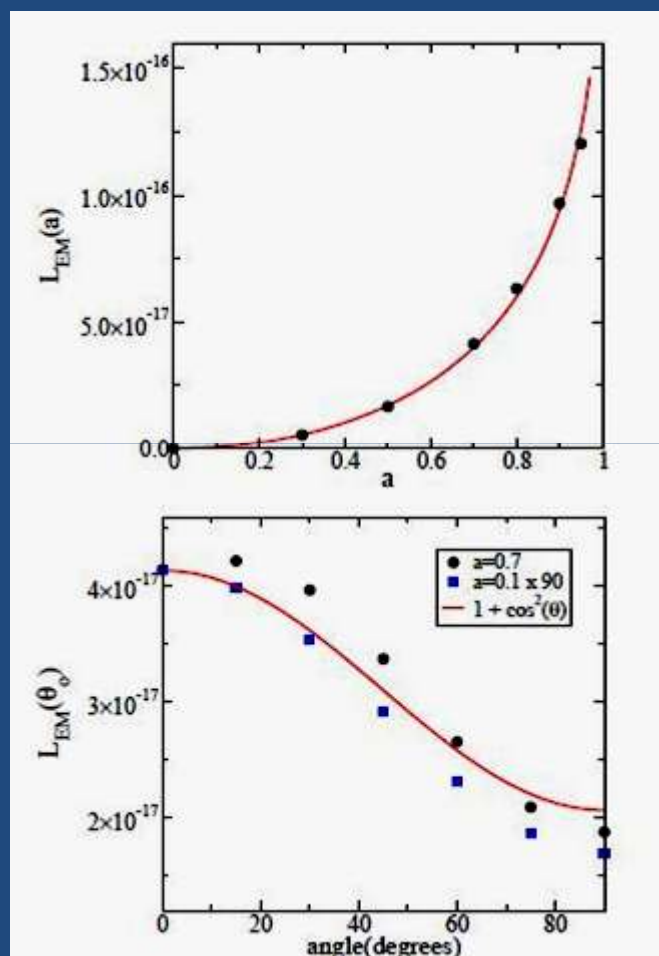
Membrane paradigm: wrt asymptotic observers, circuit moves through a B field \rightarrow EMF produced.

BH becomes the battery.

[Damour,Phinney,Thorne,McDonald...]



Single BHs, disk alignment?



- we knew. $P \sim B^2 a^2$ in the aligned case [Tchechovskoy, Narayan, McKinney 2010].
- For misaligned case?
 - Poynting flux still there, along B
 - $P \sim B^2 a^2 (1 + \cos^2)$

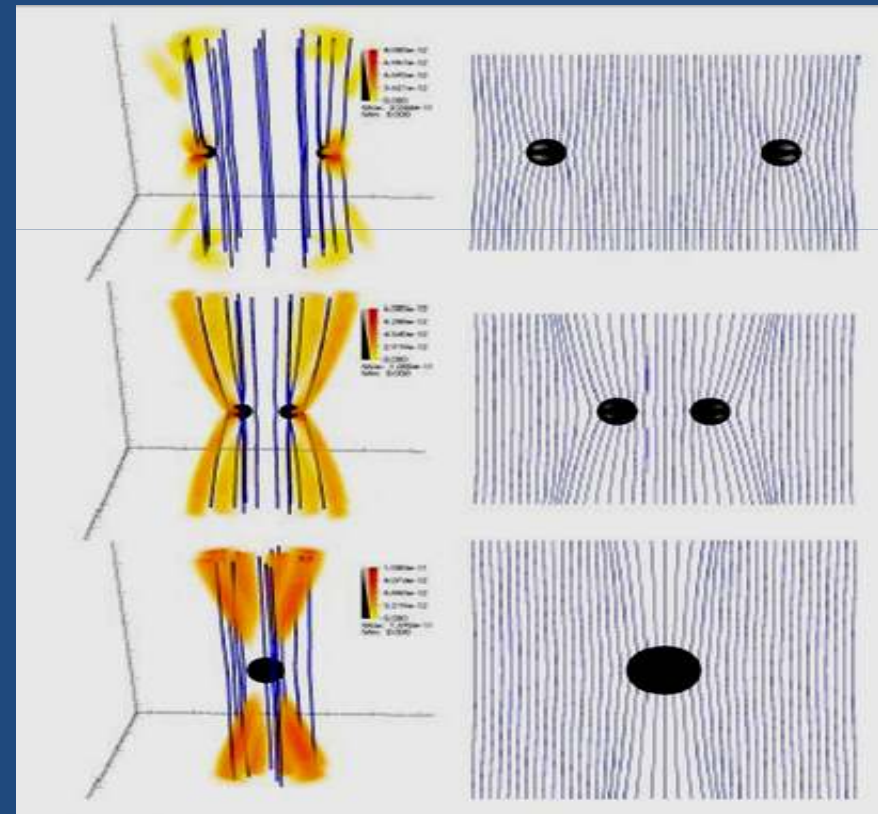
[Palenzuela, Garret, LL. Liebling, to be subm]

Onto binaries

- Head on & quasicircular, equal mass. *non-spinning*
- Magnetic field as given by a 'circular loop' at far distances \sim constant within computational domain
- Field strength $\sim 10^4\text{G}$
 - For this value, if $M_T=10^8 M_O$, EM Energy dens $\sim 10^{-16} [1/\text{M}^2]$
 - EM fields won't affect binary dynamics, but the other way around

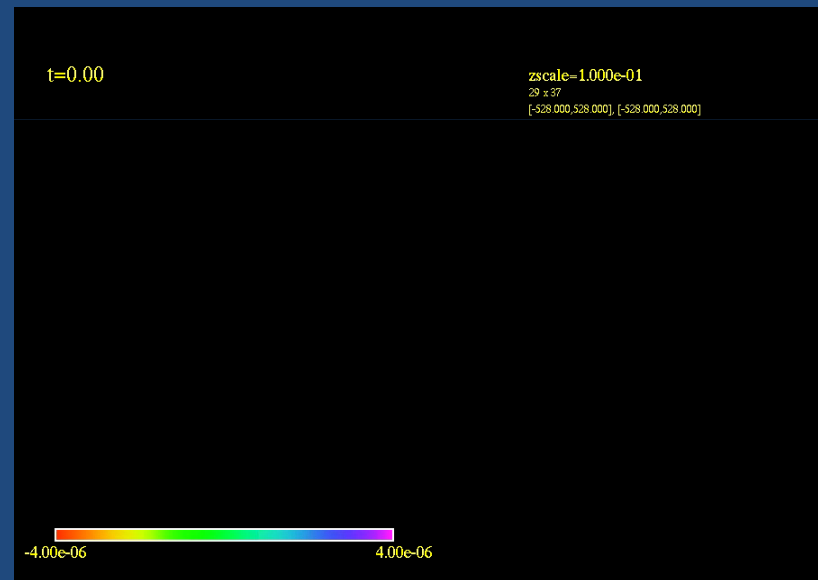
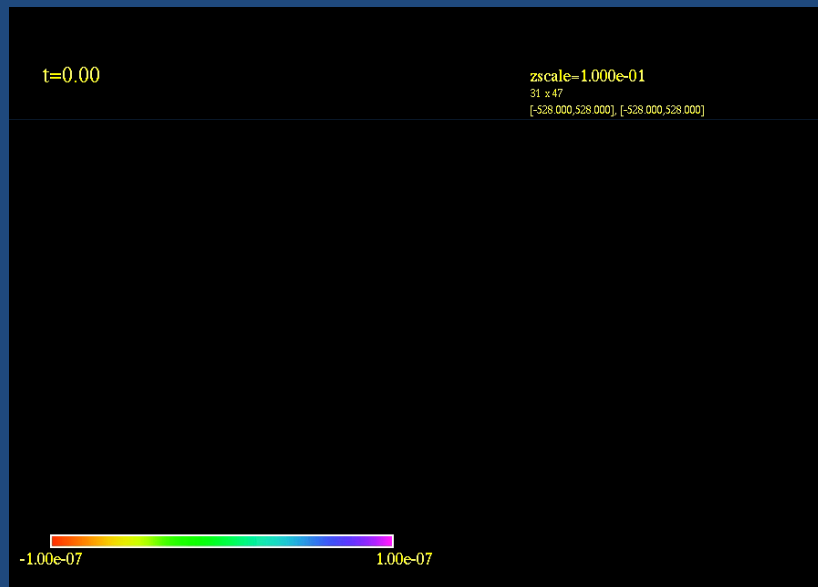
Head-on case.

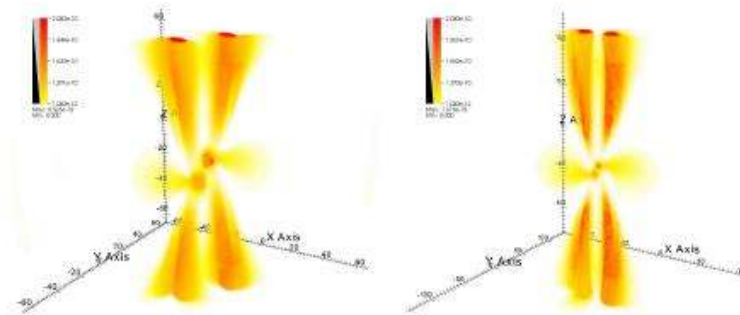
- Poynting flux,
 - What sources it ?
 - field lines tension/breaking as BH pulls them
- Membrane paradigm:
 - “Charge” separation induced by “Hall effect” , thus circuit is still there and still moving through B.
- Poynting flux induced, though shuts off after merger



Onto the binary case

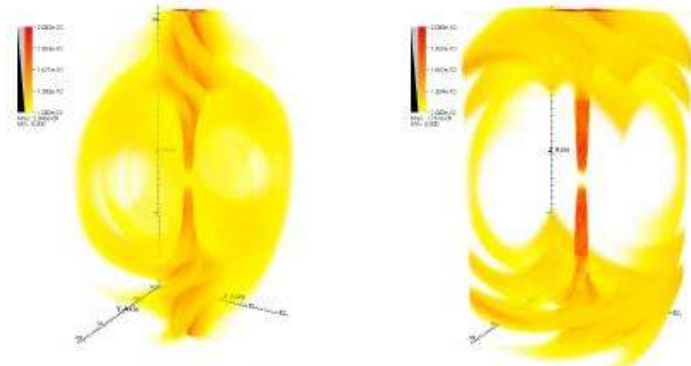
- Orbit \rightarrow Black holes move through B. As in head-on case, 'circuit' can be Established due to charge separation (see in vacuum case already, [Palenzuela et.al. Moesta et.al.])
- Thus, expect Poynting flux through orbiting stages. Also at late time (BZ).





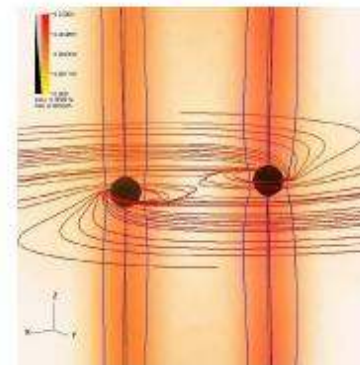
(a) $-11.0 M_8$ hrs

(b) $-3.0 M_8$ hrs

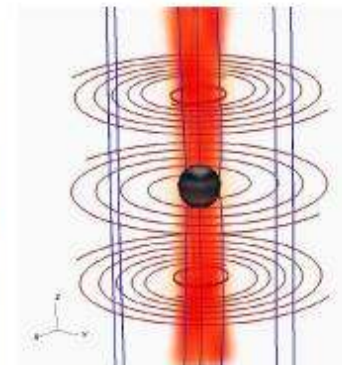


(c) $4.6 M_8$ hrs

(d) $6.8 M_8$ hrs



(a) $-8.2 M_8$ hrs

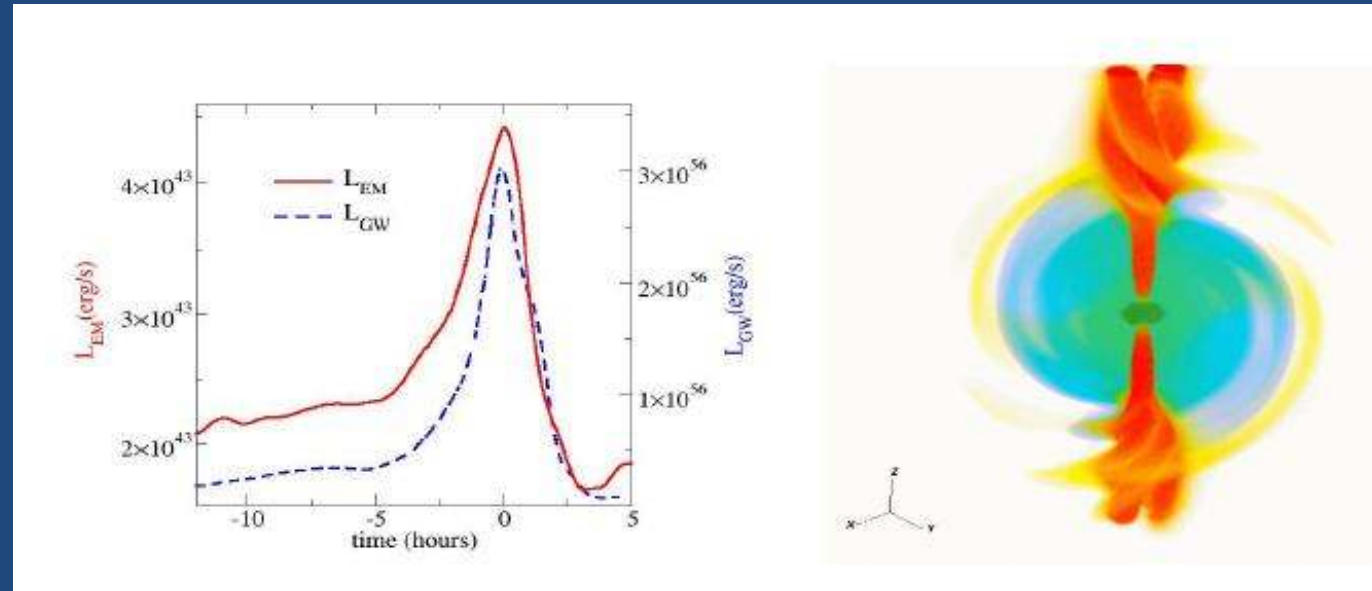


(b) $4.6 M_8$ hrs

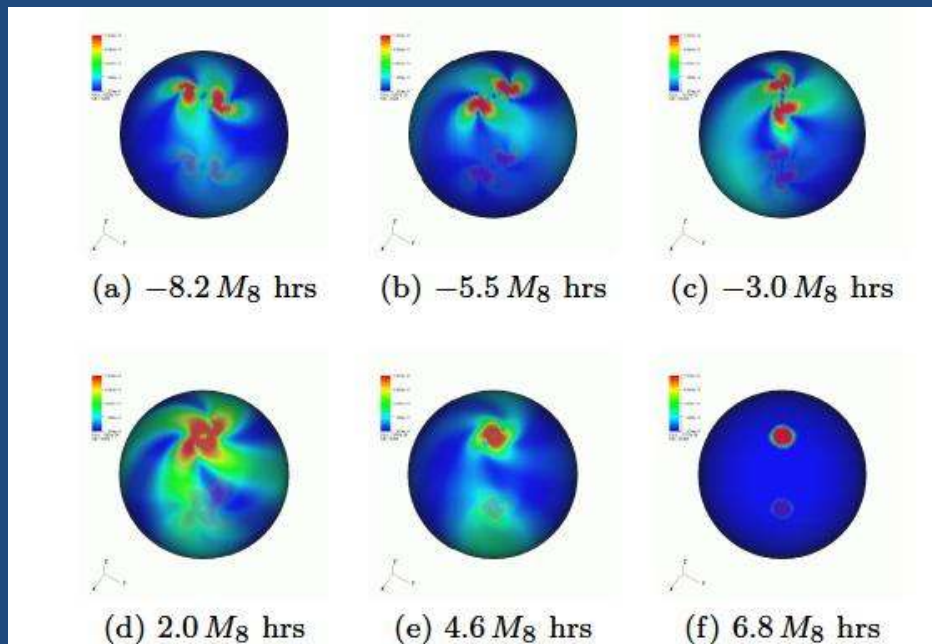
Ω_F this need not make sense!

Poynting flux

- Energy flux:



- Strong emission throughout. Burst around merger epoch



$m=2 \rightarrow 0$ transition

Distributed energy output

- Making contact with astro... recall $(R_{\text{orb}} \Omega_{\text{orb}}) < 1$
- GW energy flux $\sim R_{\text{orb}}^4 \Omega_{\text{orb}}^6 M^2$ --strong emission--
- EM energy flux $\sim (R_{\text{orb}} \Omega_{\text{orb}})^2 B^2$ --weaker but sustained, doesn't shut off after merger--
- Spinning case will have BZ on top. Also, particularly 'cute' scenarios should show an interesting phenomenology
- For 10^4G , $10^8 M_{\odot}$ flux $\sim 10^{43-44}$ ergs. IF Poynting flux energy efficiently transferred to observable emissions, interesting pre/post merger observations possible; to $z=1$?