# Dual Jets from binary BHs

CITA@25 May 14, 2010



Carlos Palenzuela<sup>(1)</sup>, L. Lehner<sup>(2)</sup>, T. Garret<sup>(2)</sup>, S. Liebling<sup>(3)</sup> (1) CITA (2) PI, UoG (3) Long Island Univ.

### Motivation : merger of galaxies

-observations indicate the presence of supermassive BHs in the center of galaxies, surrounded by gas and an accretion disk
- in the Active Galactic Nuclei (AGN), the BHs are surrounded by a disc of matter likely magnetized. For a M=10<sup>8</sup>M<sub>o</sub>

- \* bounded by the jets :  $B_0 < 10^4 10^6 G$  (near the BH)
- \* Eddington magnetic field :  $B_0 \sim 10^5 \text{ G}$



### Motivation : merger of galaxies

- the galaxies have undergone some mergers
- during the merger, the binary BH hollows
   the surrounding gas while their orbit shrinks,
   forming a circumbinary disk
- (Milosavljevic & Phinney, Astrophys. J. 622)
- eventually, the dynamics of the binary is dominated by GW





### Motivation : merger of galaxies

- the luminosity of the disk is modified by the binary BH dynamics
- the merger can enhanced some Blandford-Znajeck mechanism
- study the correlations between GW & EM radiation
- study systems with both bands to extract more information



General Relativity for the evolution of the spacetime
Maxwell equations for the evolution of the EM fields
Hydrodynamics for the evolution of the plasma
Radiation processes due to the accretion, disk dynamic..

### Motivation : pre-merger effects

before/during the merger (CP et al., 2010):
\* study the effects of the binary BHs dynamics in the EM fields and the radiation produced at the merger



- sub-domain with the BHs, excluding the disk General Relativity for the evolution of the spacetime
Maxwell equations for the evolution of the EM fields
Force-free approximation for the plasma
Radiation processes due to the accretion, disk dynamic...

### The Einstein-Maxwell system

• Einstein equations (see Abdul's and Harald's talks)

$$\mathbf{R}_{\mathbf{b}} - \mathbf{R} \, \mathbf{g}_{\mathbf{b}} / 2 = 0 \qquad \mathbf{R}_{\mathbf{b}} = \mathbf{R}_{\mathbf{b}} \left( \partial_{\mathbf{d}} \mathbf{g}_{\mathbf{b}}, \partial_{\mathbf{c}} \mathbf{g}_{\mathbf{b}}, \mathbf{g}_{\mathbf{b}} \right)$$

• Maxwell equations, with F the Maxwell tensor I the current 4-vector

$$\nabla_{a} F^{ab} = -I^{b} \qquad F^{ab} = n^{a} E^{b} - n^{b} E^{a} + \varepsilon^{ab} B_{c}$$
$$\nabla_{a} *F^{ab} = 0 \qquad \partial_{t} E - \nabla x B = -J, \quad \nabla \cdot E = q$$
$$\partial_{t} B + \nabla x E = 0, \quad \nabla \cdot B = 0$$

### The force-free approximation

If B and the black hole spin are large enough (B ≈ 0.2 T for M≈10<sup>8</sup> M<sub>☉</sub>, a>0.7), the induced E will accelerate charged particle to high enough energies that will radiate photons decaying into electron-positron pairs.
The time average of this cascade process will be a

magnetosphere full of a tenuous plasma, which can support charges and currents but with negligible inertia

$$\mathbf{\nabla}_{\mathbf{a}} \mathbf{T}^{\mathbf{a}} = \mathbf{0} \quad \Rightarrow \quad \mathbf{\nabla}_{\mathbf{a}} \mathbf{T}^{\mathbf{a}}_{(\text{fluid})} = -\mathbf{\nabla}_{\mathbf{a}} \mathbf{T}^{\mathbf{a}}_{(\text{em})} = -\mathbf{F}^{\mathbf{a}} \mathbf{I}_{\mathbf{a}}$$
  
if  $\rho, \mathbf{P} \ll \mathbf{B}^2$  then  $\mathbf{\nabla}_{\mathbf{a}} \mathbf{T}^{\mathbf{a}}_{(\text{fluid})} \ll \mathbf{F}^{\mathbf{a}} \mathbf{J}_{\mathbf{a}} \approx \mathbf{0}$ 

 $\mathbf{E} \cdot \mathbf{J} = 0$ ,  $\mathbf{q} \cdot \mathbf{E} + \mathbf{J} \cdot \mathbf{x} \cdot \mathbf{B} = 0 \rightarrow \mathbf{E} \cdot \mathbf{B} = 0$ 

### The Blandford-Znajek mechanism

• By using these conditions, we can compute the electromagnetic energy flux

$$F_E = 2 (B^r)^2 \Omega_F r \left(\frac{a}{2Mr} - \Omega_F\right) sin^2 \theta$$
$$- B^r B^\phi \Omega_F \Delta sin^2 \theta.$$

• at the AH,  $\Delta = 0$  and we define the rotation frequency of the BH like  $\Omega_{\rm H} = a/(2*M*r_{\rm H})$ 

$$F_E|_{r=r_H} = 2 \, (B^r)^2 \, \Omega_F \, r_H \, (\Omega_H - \Omega_F) \, sin^2 \theta.$$

so if B<sup>r</sup>>0 and  $0 < \Omega_{F} < \Omega_{H}$ , there will be an outward energy flux at the AH

## Single black holes

• There is a rotation of the EM field lines and an extraction of energy  $\rightarrow$  Blandford-Znajek mechanism



a = 0.99, angle =0

a = 0.99, angle = 45°

 $M = 10^8 M_{\odot}$ ,  $B = 10^4 G$ 

### Single black holes

• The radiated power as a function of

#### 8×10<sup>-5</sup> $2.5 \times 10^{-5}$ 6×10<sup>-5</sup> $\boldsymbol{L}_{\text{EM}}$ <sup>™</sup> 2.0×10<sup>-5</sup> 4×10<sup>-5</sup> 2×10<sup>-5</sup> $1.5 \times 10^{-5}$ 0 $1.0 \times 10^{-5}$ 20 80 40 60 0.2 0.4 0.6 0.8 0 angle(degrees) a

the spin

### the inclination angle

## Binary black holes : head on

### • The EM radiation propagates along the magnetic field lines



## **Binary** black holes : inspiral-merger

### The EM radiation is collimated !!



## Binary black holes : inspiral-merger

### Transition for m=2 to m=0 (decays to Blandford-Znajek)



### **Binary** black holes : inspiral-merger

• The EM power ~  $(B\Omega R)^2$ , while the GW power ~  $M^2 R^4 \Omega^6$ .

• A significant amount of EM energy is radiated days/weeks before the merger, while most of the GW is emitted during the last day



## Summary

- we have evolved EE + Maxwell + force-free to study the effects of the binary BH dynamics on the EM fields

- EM energy radiation even if the BHs are not spinning.

- Persistent collimated EM radiation, enhanced at the merger

- Future steps :

- \* consider spinning cases
- \* go beyond the force-free approximation
- \* include a realistic disk model
- \* include radiation....